



Beijing, China

# Numerical simulation on gas combustion and spray flames

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## Beihang University



- It was established in 1952. Last year, it is the 60 anniversary !
- It comprises 27 schools, 17 academician and 3759 faculties. The enrollment is over 27800, including 14428 UG, 4015 Ph.D candidates, 668 oversea.

## Ranking and the School

- ✿ General ranking: 7-15 in China;
- ✿ Aerospace propulsion Theory and Engineering, No.1, Power Engineering and Engineering Thermal Physics, No.7.
- ✿ Our school has **four** departments and **one** national key laboratory: Department of Aviation propulsion, Department of Fluid Machinery, Department of Engineering Thermophysics, Department of Thermal Engineering, The National Key Laboratory on Aero-Engines.

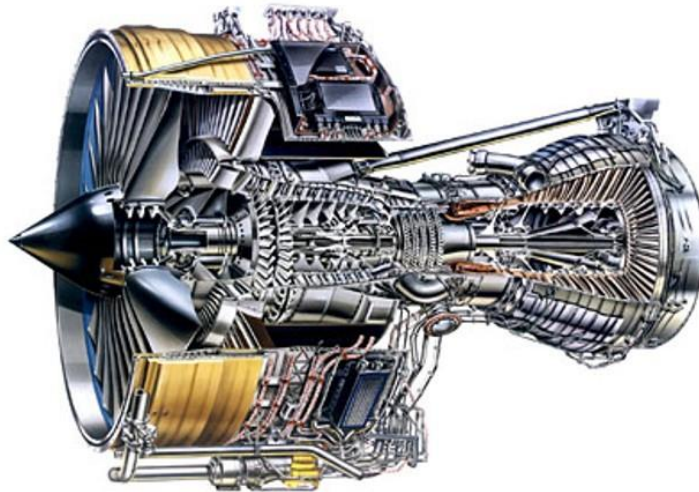


## Content

- Numerical simulation on
  - gas combustion : An algebraic sub-grid scale turbulent combustion model
  - spray flames : An droplet burning model



## **Gas turbine combustor**



- **Gas-liquid two phase turbulent combustion**
- **Droplets tracing, chemical kinetics, turbulent flow in complex geometry, turbulent reaction rate, radiation, ...**

## Turbulent combustion rate

- RANS, LES, DNS are the three strategies...
- Turbulent combustion rate is a big problem in combustion simulation, because of the exponential part.

$$w_i = B\rho^2 Y_1 Y_2 \exp\left(-\frac{E}{RT}\right)$$

- The average reaction rate is not the function of the averaged value
- Couple with turbulence...another big problem
- Turbulent combustion models are needed for engineering applications.

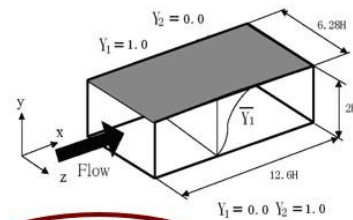
## Background – about SGS combustion model

- Recently, **Large eddy simulation (LES)** is becoming used in engineering applications.
- For **gas turbine combustion chambers**, it is used to solve the flow topology, emission, compressor/turbine interaction, cooling, ...problems
- The **sub-grid scale (SGS) combustion model** is one of the key aspects of LES method.
- There are various SGS combustion models: the *laminar flamelet* model (Pitsch et al., Senoner et al. and Mattsson et al.), *assumed PDF* combustion model (Moin et al, James et al.), *linear eddy* model (Menon et al.), *EBU* model (Menon et al.), *PDF equation* model (James et al.), *Dynamically Thickened Flame* combustion model (Poinsot et al. ) ...
- The SGS turbulent combustion model is still in developing

## Background - DNS verification and LES validation

- As an important tool for fundamental studies, **direct numerical simulation (DNS)** attracted increasing interest in recent years.
- Attention is paid to finding detailed instantaneous flow and flame structures, understanding the turbulence/flame interaction (Poinsot et al. , Zhang and Rutland) and verifying RANS and LES models (Zhang and Rutland, Luo, Overholt and Pope, Bedat et al., Sreedhara and Huh, Hawkes and Chen, Chong and Heinz).
- So the **SGS model** should be better **verified by DNS** before applying to test cases. Thus, in this paper:
  - A direct numerical simulation of turbulent reacting channel flows, taking the buoyancy effects account, is performed using a spectral method, and the **DNS database** is used to **verify the SGS combustion model**
  - This SGS combustion model is tested using methane and **methanol-air jet flame**.

## DNS Governing Equations



$$\frac{\partial u_i}{\partial x_i} = 0$$

$$\frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j} = -\frac{\partial p}{\partial x_i} + \frac{\partial}{\partial x_j} \left[ \nu \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \right] - F_i \quad F_y = -g \frac{T - T_w}{T_w}$$

$$\frac{\partial}{\partial t} (Y_i) + \frac{\partial}{\partial x_j} (u_j Y_i) = \frac{\partial}{\partial x_j} \left( D \frac{\partial Y_i}{\partial x_j} \right) - w_i / \rho \quad i = 1, 2$$

The buoyancy effect using Boussinesq approximation

$$\frac{\partial}{\partial t} (T) + \frac{\partial}{\partial x_j} (u_j T) = \frac{\partial}{\partial x_j} \left( \frac{\lambda}{c_p \rho} \frac{\partial T}{\partial x_j} \right) + w_1 Q_1 / (c_p \rho)$$

One-step Arrhenius chemical kinetics

$$w_i = B \rho^2 Y_1 Y_2 \exp\left(-\frac{E}{RT}\right)$$

## Numerical Methods

- Galerkin-Tau spectral expansion method: Fourier transform is used in x and z directions and the Chebyshev transform is used in the y direction;
- Uniform grid distribution is used in x and z directions and the Gauss-Lobatto nonuniform grid distribution is used in the y direction.  $128 \times 128 \times 129$  results in a total of 2.11 million nodes. The time step used is  $0.01H/U_m$  with a third-order scheme.
- For all cases the mass fraction of species 1 (fuel) is given as 1.0 at the top wall and 0.0 at the bottom wall, whereas the mass fraction of species 2 (oxidizer) is given as 0.0 at the top wall and 1.0 at the bottom wall. The wall temperature is given as 900K.
- Periodic boundary conditions are used in the longitudinal and span-wise directions and solid-wall boundary conditions are used on the top and bottom boundaries.

### Parameters for DNS Cases

$$F_y = -g \frac{T - T_w}{T_w}$$

Case	B	E/R (K)	Q (kJ/kg)	$F_y$ (m/s <sup>2</sup> )
1	0	0	0	0
2	0.1	0	0	0
3	1.0	0	0	0
4	$10^8$	15000	100	0
5	$10^8$	20000	100	0
6	$10^{10}$	20000	100	0
7	$10^8$	20000	100	-0.5
8	$10^8$	15000	100	$F_y$
9	$10^8$	20000	100	$F_y$

$$\frac{\partial}{\partial t}(T) + \frac{\partial}{\partial x_j}(u_j T) = \frac{\partial}{\partial x_j} \left( \frac{\lambda}{c_p \rho} \frac{\partial T}{\partial x_j} \right) + w_1 Q_1 / (c_p \rho)$$

$$w_i = B \rho^2 Y_1 Y_2 \exp\left(-\frac{E}{RT}\right) = \rho^2 K Y_1 Y_2$$

# LES Governing Equations

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_i} (\rho \tilde{u}_i) = 0$$

$$\frac{\partial}{\partial t} (\rho \tilde{u}_i) + \frac{\partial}{\partial x_j} (\rho \tilde{u}_i \tilde{u}_j) = \frac{\partial}{\partial x_j} \left( \mu \frac{\partial \tilde{u}_i}{\partial x_j} \right) - \frac{\partial \tilde{P}}{\partial x_i} - \frac{\partial \tau_{ij}}{\partial x_j}$$

**Smagorinsky-Lilly's eddy viscosity model**

$$\tau_{ij} = -2\mu_t \tilde{S}_{ij} + \frac{1}{3} \tau_{kk} \delta_{ij}$$

$$\frac{\partial \rho \tilde{Y}_s}{\partial t} + \frac{\partial}{\partial x_j} (\rho \tilde{u}_j \tilde{Y}_s) = \frac{\partial}{\partial x_j} \left( \frac{\mu}{S_{cs}} \frac{\partial \tilde{Y}_s}{\partial x_j} \right) - \tilde{w}_s - w_{s,s} - \frac{\partial g_{sj,s}}{\partial x_j}$$

$$\frac{\partial \rho \tilde{h}}{\partial t} + \frac{\partial}{\partial x_j} (\rho \tilde{u}_j \tilde{h}) = \frac{\partial}{\partial x_j} \left( \frac{\mu}{Pr} \frac{\partial \tilde{h}}{\partial x_j} \right) - \frac{\partial q_{sj}}{\partial x_j}$$

**The filtered reaction rate and the SGS reaction rate**

$$g_{sj,s} = \frac{\mu_t}{Sc_{t,s}} \frac{\partial \tilde{Y}_s}{\partial x_j} \quad q_{sj} = \frac{\mu_t}{Pr_t} \frac{\partial \tilde{T}}{\partial x_j}$$

## ASSCM SGS Combustion Model

$$\frac{\partial \rho \tilde{Y}_s}{\partial t} + \frac{\partial}{\partial x_j} (\rho \tilde{u}_j \tilde{Y}_s) = \frac{\partial}{\partial x_j} \left( \frac{\mu}{S_{cs}} \frac{\partial \tilde{Y}_s}{\partial x_j} \right) - \tilde{w}_s - w_{s,s} - \frac{\partial g_{sj,s}}{\partial x_j}$$

**The filtered reaction rate**

**The SGS reaction rate**

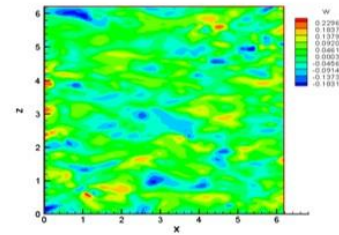
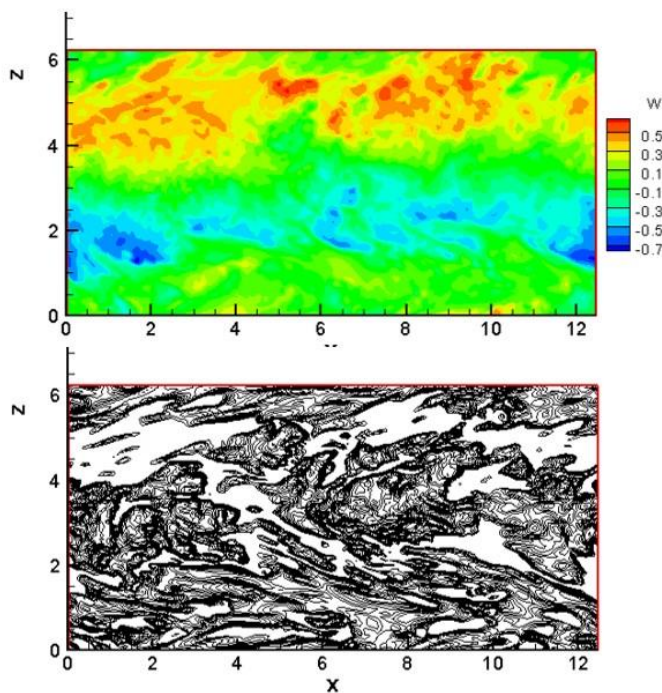
$$\tilde{w}_s = \rho^2 \tilde{K} \tilde{Y}_{OX} \tilde{Y}_{Fu}$$

$$w_{s,s} = \rho^2 [K (\tilde{Y}_{OX} \tilde{Y}_{Fu} - \tilde{Y}_{OX} \tilde{Y}_{Fu}) + Y_{OX} (\tilde{K} \tilde{Y}_{Fu} - \tilde{K} \tilde{Y}_{Fu}) + Y_{Fu} (\tilde{K} \tilde{Y}_{OX} - \tilde{K} \tilde{Y}_{OX})]$$

$$(\tilde{K} \tilde{Y}_{OX} - \tilde{K} \tilde{Y}_{OX}) = C_{K,Y_{OX}} L_s^2 \frac{\partial \tilde{K}}{\partial x_j} \frac{\partial \tilde{Y}_{OX}}{\partial x_j}$$

**The constant C is 0.005**

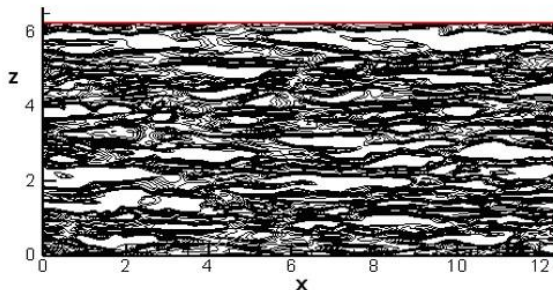
## The DNS Instantaneous Results (1)



**Instantaneous  
Longitudinal Velocity  
(case 8)**

**Instantaneous  
Concentration Fluctuation  
(case 8)**

## The DNS Instantaneous Results (2)



**Instantaneous  
Concentration  
Fluctuation (case 5)**

- In cases without scalar-velocity interaction, the velocity is rather independent of the scalar, and the scalar transportation does not affect the velocity field. Typically the velocity/scalar field is full of the strip structures.
- In the interaction cases (case 8), The buoyancy term is a moderate two-way coupling between the velocity and the scalar transportation. The buoyancy force is in the y direction, but the transportation induced by it can be seen in the longitudinal (z) direction.



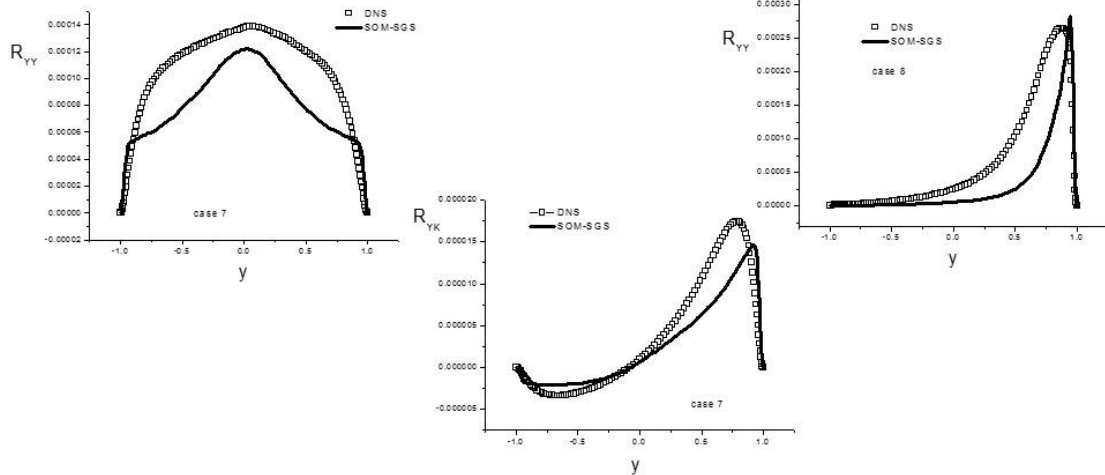
# The DNS Verification (1)

The ASSCM model is generally given as:

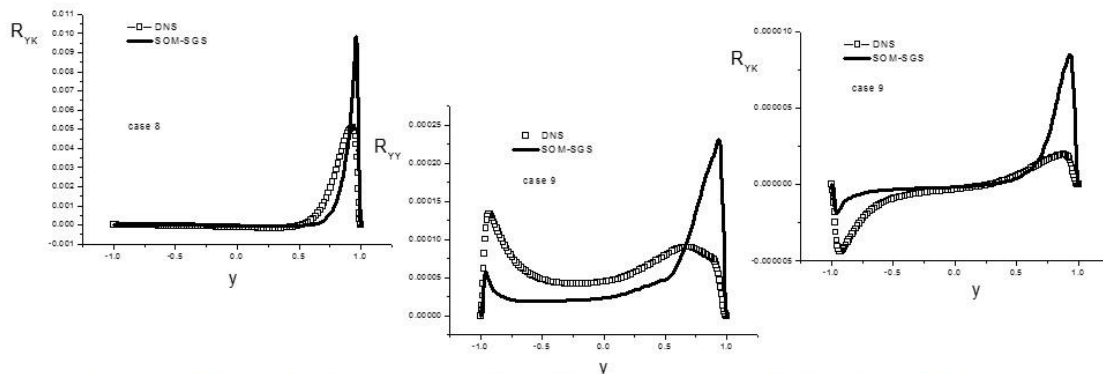
$$R_{\Phi\Psi} = \underbrace{\widetilde{\Phi\Psi} - \widetilde{\Phi}\widetilde{\Psi}}_{\text{Model}} - C_{\Phi,\Psi} L_s^2 \frac{\partial \widetilde{\Phi}}{\partial x_j} \frac{\partial \widetilde{\Psi}}{\partial x_j}$$

Directly from instantaneous DNS database

From the filter-averaged DNS database



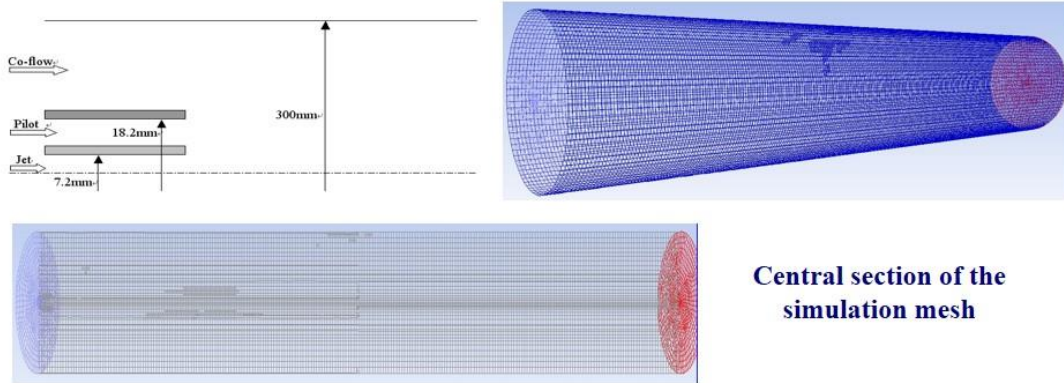
# DNS Verification (2)



- The model results has the same trend as the exact solution. For all 3 cases, the  $R_{YY}$  value remains in the same range, while the  $R_{KY}$  value changes significantly
- There are maximum errors of about 30%, 50% and even over 100%...
- Though these cases are not ideal examples for real reaction flow, they give a theoretical understanding of the model. In some regions, the ASSCM combustion model is close to the exact solution, so further validation should be carried out.

# LES Validation - Gas Jet Flame (Flame D)

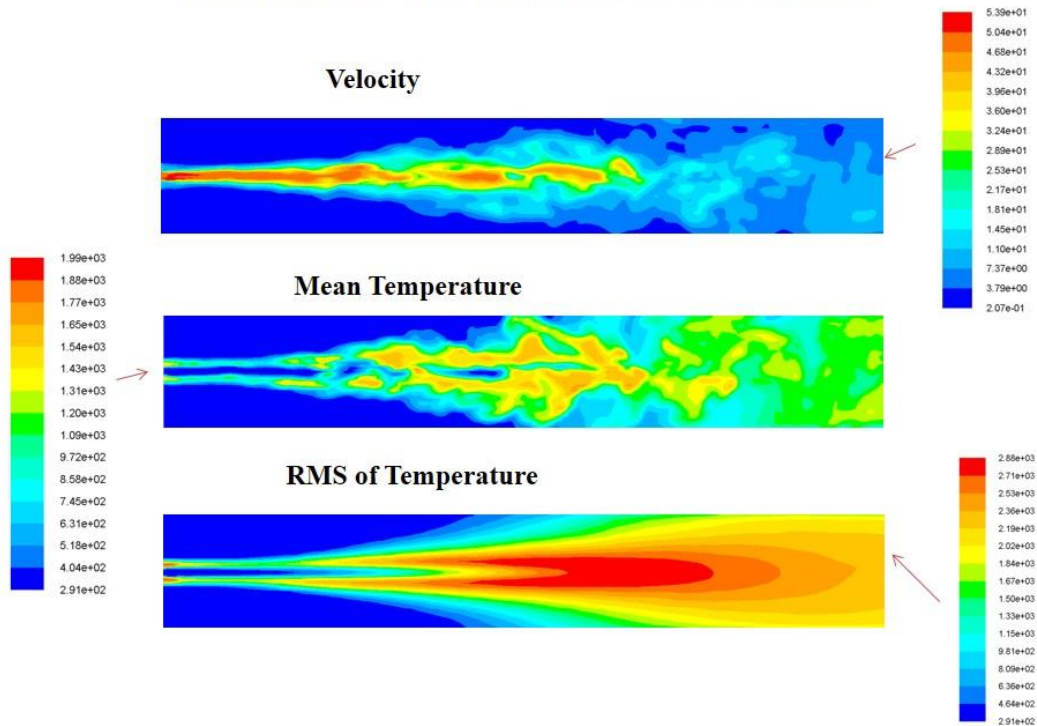
## Jet flame structure and the simulation mesh



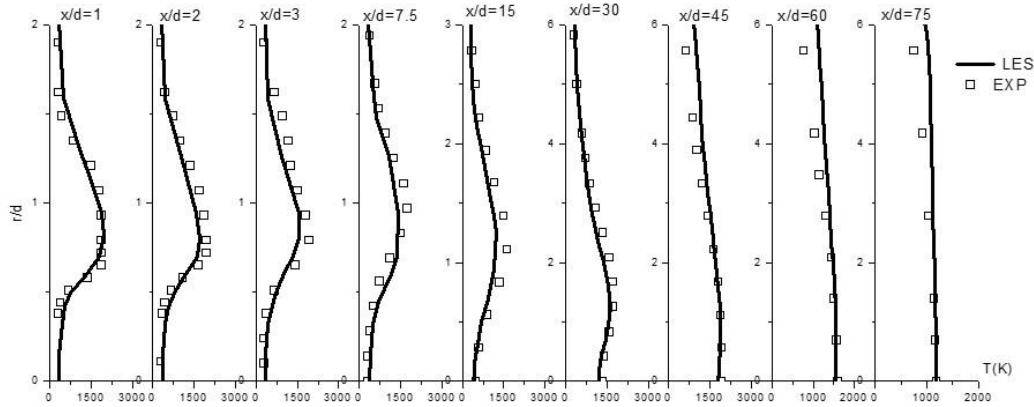
### The methane-air reaction kinetics :

$$w_{fu} = 2.119 \times 10^{11} Y_{ox}^{1.3} Y_{fu}^{0.2} \exp(-2.027 \times 10^8 / RT)$$

## Instantaneous Contours for Flame D



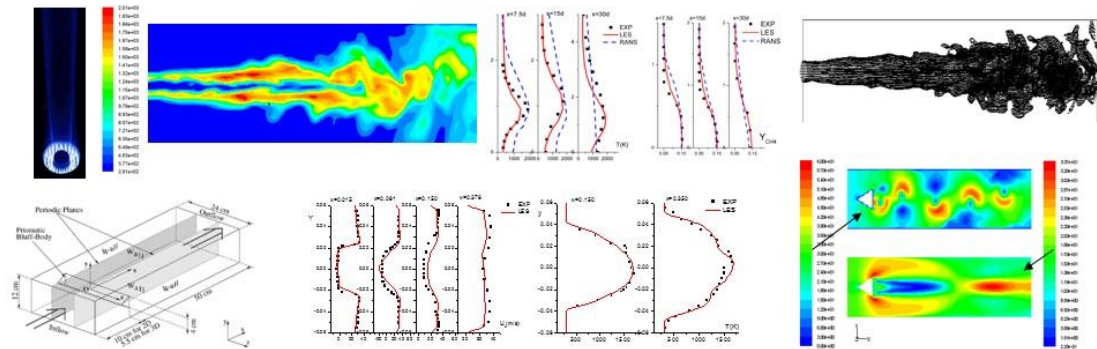
# The Average Temperature



The error is less than 5%

## ASSCM SGS model validation

- We have proposed *the algebraic sub-grid turbulent combustion model (ASSCM)* and used it in premixed combustion and a partly-diffusion flame. The statistical results were close to measurements.
- Here, we applied it for *two-phase* combustion cases.

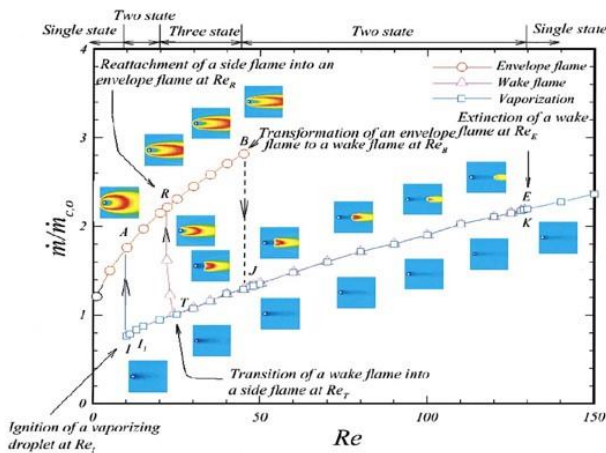


F. Wang, L. X. Zhou, C. X. Xu, Large-Eddy Simulation of Premixed Combustion and Validation of an algebraic second order moment RANS Combustion Model, *Journal of Propulsion Technology*, 2008, 29(1): 33-36  
 F. Wang, L. X. Zhou, C. X. Xu, Large-eddy simulation of correlation moments in turbulent combustion and validation of the RANS-SOM combustion model, *Fuel*, 2006, 85(9): 1242-1247

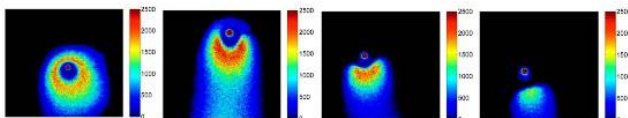
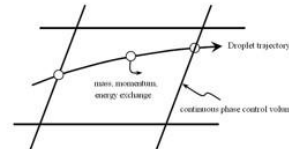
# Content

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## Full Status two phase combustion model



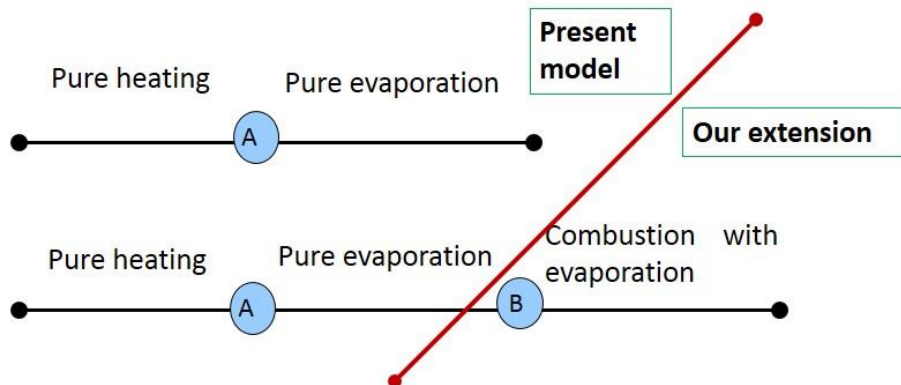
There is single droplet burning in combustion chamber and spray flames  
Need to be modeling



H H Chiu, AIAA 95-2427, A new droplet model for spray combustion, 1995

X. Mercier, M. Orain, F. Grisch, Appl. Phys. B 88, 151–160 (2007)

## The full status droplet combustion model (1)



The upper one is the general treatment of droplet, while the lower one is our new model;

The critical point B is from the droplet ignition studies; and the correlation of the droplet combustion with evaporation state is from single droplet combustion studies.

## The full status droplet combustion model (2)

- For the criterion for point B, three judgments are used:
  1. The droplet diameter is bigger than 0.001mm;
  2. The ambient temperature is higher than 1200 K;
  3. The evaporation time is longer than the ignition time.
- The ignition delay time function is:

$$\tau_i = A \exp\left(\frac{E}{RT_\infty}\right) C_{\text{oxygen}}^{-0.5} C_{\text{methanol}}^{-0.1}$$

- The evaporation time is defined as:  $\tau_e = d_{d0}^2 / k_e$
- As for the burning correlation, the traditional D<sup>2</sup> law is taken, here the  $K_c$  is 0.7 mm<sup>2</sup>/s.

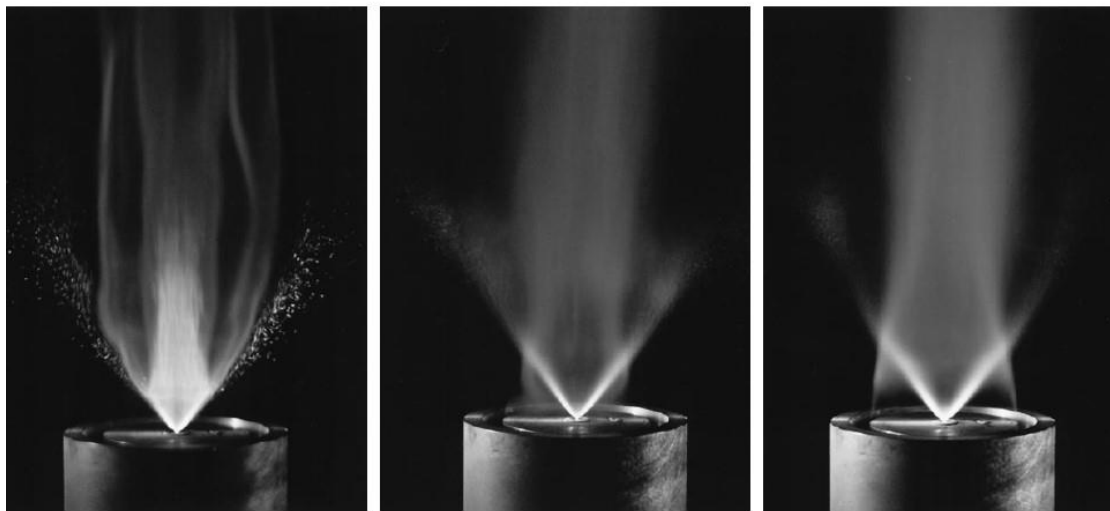
$$k_c = -d(d_d^2) / dt$$

$$\frac{dT_d}{dt} = \frac{Nu C p_g}{3 Pr_g C p_l} \left( \frac{\bar{T}_g - T_d}{\tau_d} \right) + \frac{h_{fg} m_d}{C p_l m_d}$$

$$m_d = \frac{dm_d}{dt} = - \frac{m_d}{3 Sc_g \tau_d} (Sh + Sh^{st}) \ln(1 + B)$$

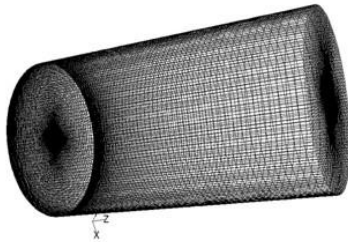
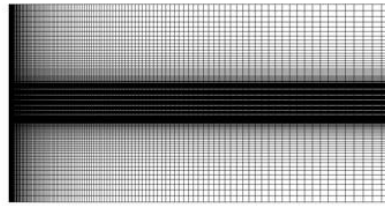
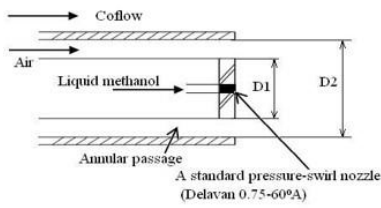
After burning:

$$\frac{d(d_p)}{dt} = \frac{4k_\infty}{\rho_p C_{p,\infty} d_p} (1 + 0.23 \sqrt{Re_d}) \ln \left[ 1 + \frac{C_{p,\infty} (T_\infty - T_p)}{h_{fg}} \right]$$



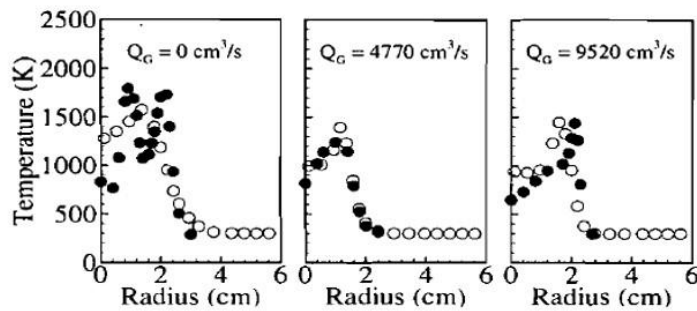
**More air, more gas mixture combustion preferential**

J. A. Friedman, M. Renksizbulut, Investigating a methanol spray flame Interacting with an annular air jet using phase-doppler interferometry and planar laser-induced fluorescence, *Combustion and Flame*, 117:661 - 684, 1999

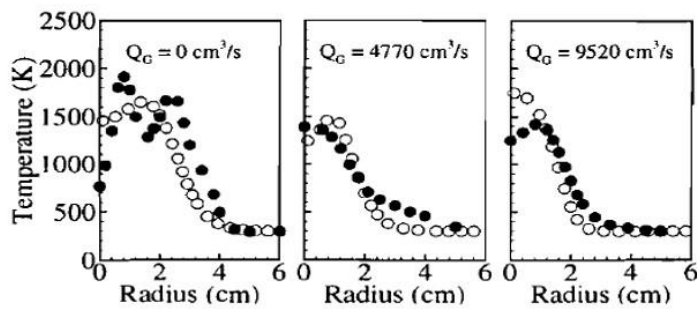


LES in gas phase;  
ASSCM turbulent  
combustion model;

$$w_{fu} = 1.799 \times 10^{10} \rho^2 Y_{CH_3OH}^{0.25} Y_{O_2}^{1.5} \exp\left(\frac{-1.256 \times 10^8}{RT}\right)$$

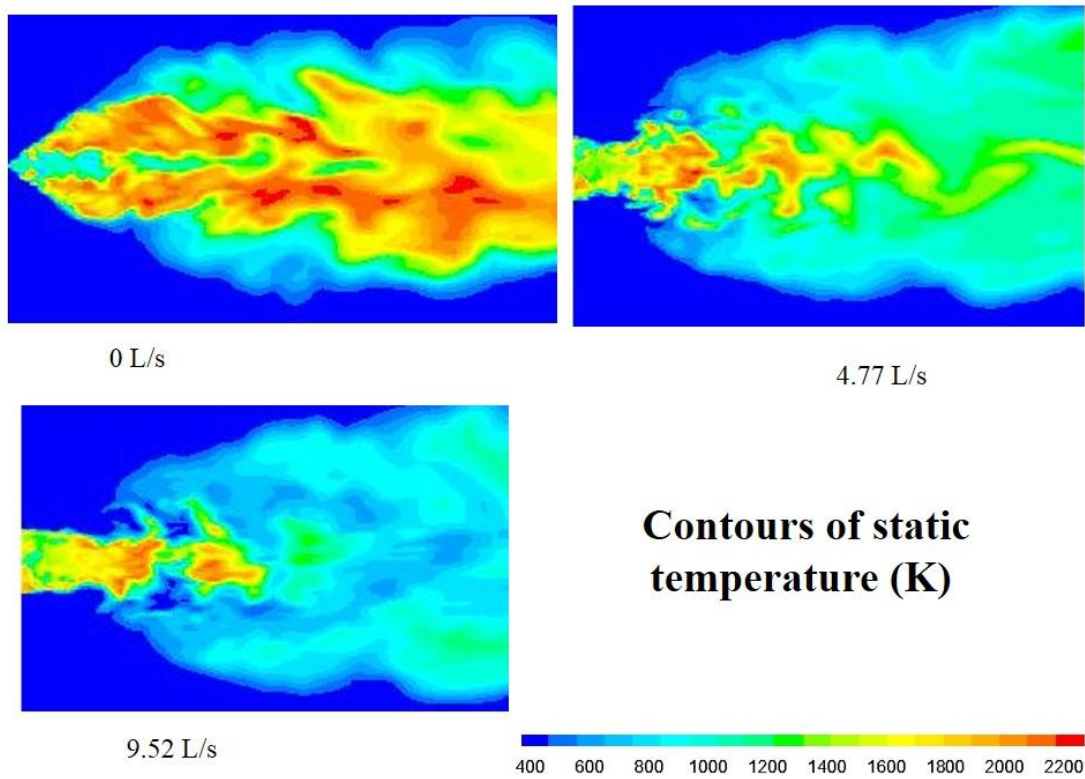


(a)  $z = 2.5$  cm

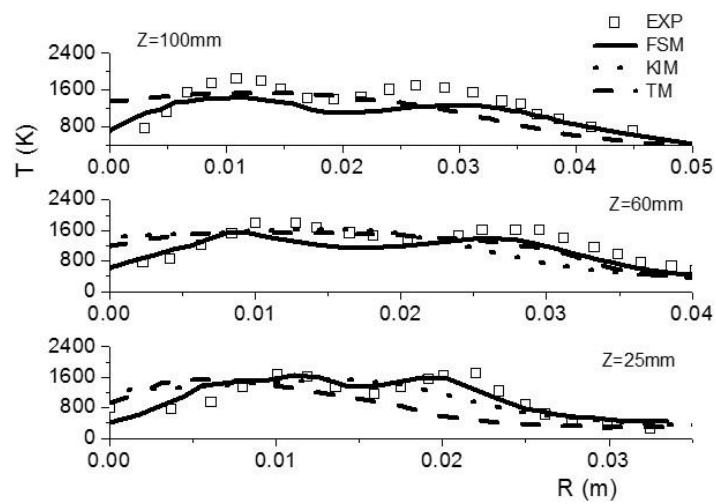


(b)  $z = 8.0$  cm

W.T. Kim, K.Y. Huh and J.A. Friedman, and M. Rensizbulut. Numerical simulation of a steady hollow-cone methanol spray flame within an annular air jet. *Combust. Sci. Technol.* 171:119-139,2001



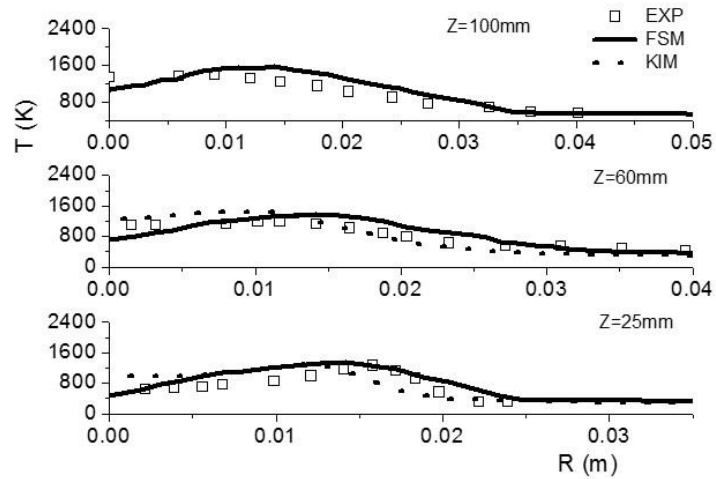
## Temperature profiles (case 1)



FSDCM - new model; Original – traditional model;  $\square$  - experimental data

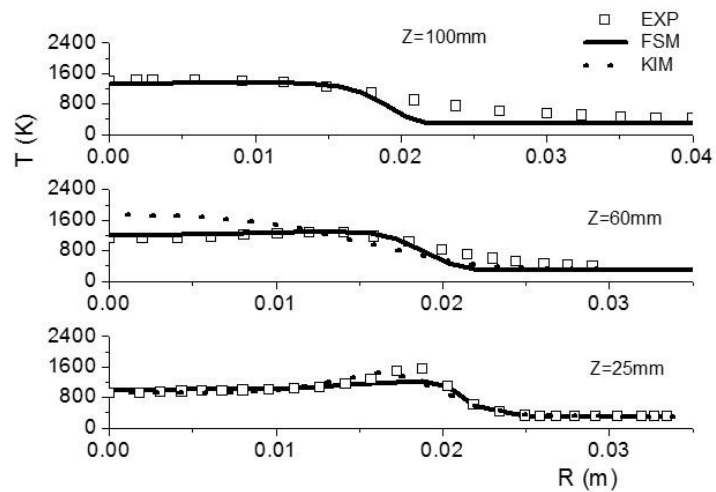


## Temperature profiles (case 2)



FSDCM - new model;  $\square$  - experimental data

## Temperature profiles (case 3)



FSDCM - new model;  $\square$  - experimental data

Thank you!

