

A model for a bi-component droplet heating and evaporation

Eng. Ahmed Elwardany

The Sir Harry Ricardo Laboratories
Centre for Automotive Engineering

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Introduction

- For a **single component** droplet, there is **only** heat diffusion in the **liquid** phase and there are **heat and mass** diffusion in the **gas** phase.
- For a **bi-component** droplet, there are **heat and mass** diffusion in both **liquid and gas** phases.

Heat diffusion in the liquid phase

- Heat conduction equation inside the droplet:

$$\frac{\partial T}{\partial t} = \kappa \left(\frac{\partial^2 T}{\partial R^2} + \frac{2}{R} \frac{\partial T}{\partial R} \right)$$

- Boundary conditions without evaporation:

$$h(T_g - T_s) = k_l \left. \frac{\partial T}{\partial R} \right|_{R=R_d-0},$$

$$\left. \frac{\partial T}{\partial R} \right|_{R=0} = 0,$$

Heat diffusion in the liquid phase

- The solution of the heat conduction equation without evaporation for $h = \text{const}$ is presented as follows:

$$T(R, t) = \frac{R_d}{R} \sum_{n=1}^{\infty} \left\{ q_n \exp \left[-\kappa_R \lambda_n^2 t \right] - \frac{\sin \lambda_n}{\|v_n\|^2 \lambda_n^2} \mu_0(0) \exp \left[-\kappa_R \lambda_n^2 t \right] - \frac{\sin \lambda_n}{\|v_n\|^2 \lambda_n^2} \int_0^t \frac{d\mu_0(\tau)}{d\tau} \exp \left[-\kappa_R \lambda_n^2 (t - \tau) \right] d\tau \right\} \sin \left[\lambda_n \left(\frac{R}{R_d} \right) \right] + T_g(t),$$

Heat diffusion in the liquid phase

- λ_n are the solutions of the equation:

$$\lambda \cos \lambda + h_0 \sin \lambda = 0,$$

$$\|v_n\|^2 = \frac{1}{2} \left(1 - \frac{\sin 2\lambda_n}{2\lambda_n} \right) = \frac{1}{2} \left(1 + \frac{h_0}{h_0^2 + \lambda_n^2} \right),$$

$$q_n = \frac{1}{R_d \|v_n\|^2} \int_0^{R_d} \tilde{T}_0(R) \sin \left[\lambda_n \left(\frac{R}{R_d} \right) \right] dR,$$

$$\kappa_R = \frac{k_l}{c_l \rho_l R_d^2}, \quad \mu_0(t) = \frac{h T_g(t) R_d}{k_l},$$

- $h_0 = (h R_d / k_l) - 1$

Heat diffusion in the liquid phase

- The effect of droplet **evaporation** has been taken into account by replacing **gas** temperature by the so-called **effective** temperature.

$$T_{\text{eff}} = T_g + \frac{\rho_l L \dot{R}_d}{h}$$

Mass diffusion in the liquid phase

- Mass fraction equation inside the droplet:

$$\frac{\partial Y_{li}}{\partial t} = D_l \left(\frac{\partial^2 Y_{li}}{\partial R^2} + \frac{2}{R} \frac{\partial Y_{li}}{\partial R} \right)$$

- Boundary conditions:

$$\alpha_i(\epsilon_i - Y_{lis}) = -D_l \left. \frac{\partial Y_{li}}{\partial R} \right|_{R=R_d-0}$$
$$\left. \frac{\partial Y_{li}}{\partial R} \right|_{R=0} = 0$$

- where $\alpha_i = \frac{D_g \rho_{\text{total}} \ln(1 + B_M)}{\rho_l R_d}$

Mass diffusion in the liquid phase

- The solution of the mass fraction equation for $\alpha_i = \text{const}$ is presented as follows:

$$Y_{li}(R, t) = \frac{R_d}{R} \sum_{n=1}^{\infty} \left\{ q_n \exp \left[-\kappa_R \lambda_n^2 t \right] - \frac{\sin \lambda_n}{\|v_n\|^2 \lambda_n^2} \mu_0(0) \exp \left[-\kappa_R \lambda_n^2 t \right] - \frac{\sin \lambda_n}{\|v_n\|^2 \lambda_n^2} \int_0^t \frac{d\mu_0(\tau)}{d\tau} \exp \left[-\kappa_R \lambda_n^2 (t - \tau) \right] d\tau \right\} \sin \left[\lambda_n \left(\frac{R}{R_d} \right) \right] + \epsilon_i(t),$$

Mass diffusion in the liquid phase

- λ_n are the solutions of the equation:

$$\lambda \cos \lambda + h_0 \sin \lambda = 0,$$

$$\|v_n\|^2 = \frac{1}{2} \left(1 - \frac{\sin 2\lambda_n}{2\lambda_n} \right) = \frac{1}{2} \left(1 + \frac{h_0}{h_0^2 + \lambda_n^2} \right),$$

$$q_n = \frac{1}{R_d \|v_n\|^2} \int_0^{R_d} \tilde{Y}_{li0}(R) \sin \left[\lambda_n \left(\frac{R}{R_d} \right) \right] dR,$$

$$\kappa_R = \frac{D_l}{R_d^2}, \quad \mu_0(t) = -\frac{\alpha_i \epsilon_i(t) R_d}{D_l},$$

- $h_0 = (-\alpha_i R_d / D_l) - 1$ and $\epsilon_i = \frac{Y_{vis}}{\sum_i Y_{vis}}$,

Heat and mass transfer

- The total evaporation rate of the droplet:

$$\dot{m} = \pi \rho_{\text{ref}} D D_i B_{M,i} Sh_i$$

- Heat transfer coefficient (h) is calculated using Nusselt number equation (Nu):

$$Nu = 2R_d h / k_g$$

Heat and mass transfer

- Nusselt number

$$Nu_{iso} = \frac{\ln(1 + B_T)}{B_T} \left(2 + \frac{0.6Re^{1/2}Pr^{1/3}}{F(B_T)} \right)$$

- Sherwood number

$$Sh_{i,iso} = \frac{\ln(1 + B_{M,i})}{B_{M,i}} \left(2 + \frac{0.6Re^{1/2}Sc_i^{1/3}}{F(B_{M,i})} \right)$$

Heat and mass transfer

- The effect of interacting between droplets was taken into account using the following equation :

$$\eta(C) = \frac{Nu}{Nu_{iso}} = \frac{Sh}{Sh_{iso}}$$

- where C is the distance parameter (distance between droplets divided by their diameters)

$$\eta(C) = 1 - 0.57 \left(1 - \frac{1 - 0.57e^{-0.13(C-6)}}{1 + 0.57e^{-0.13(C-6)}} \right)$$

Species mass fractions at the droplet surface

- Knowing the liquid mass fraction at droplet surface Y_{li} as initial value or from the solution of species diffusion equation then we can calculate the following:

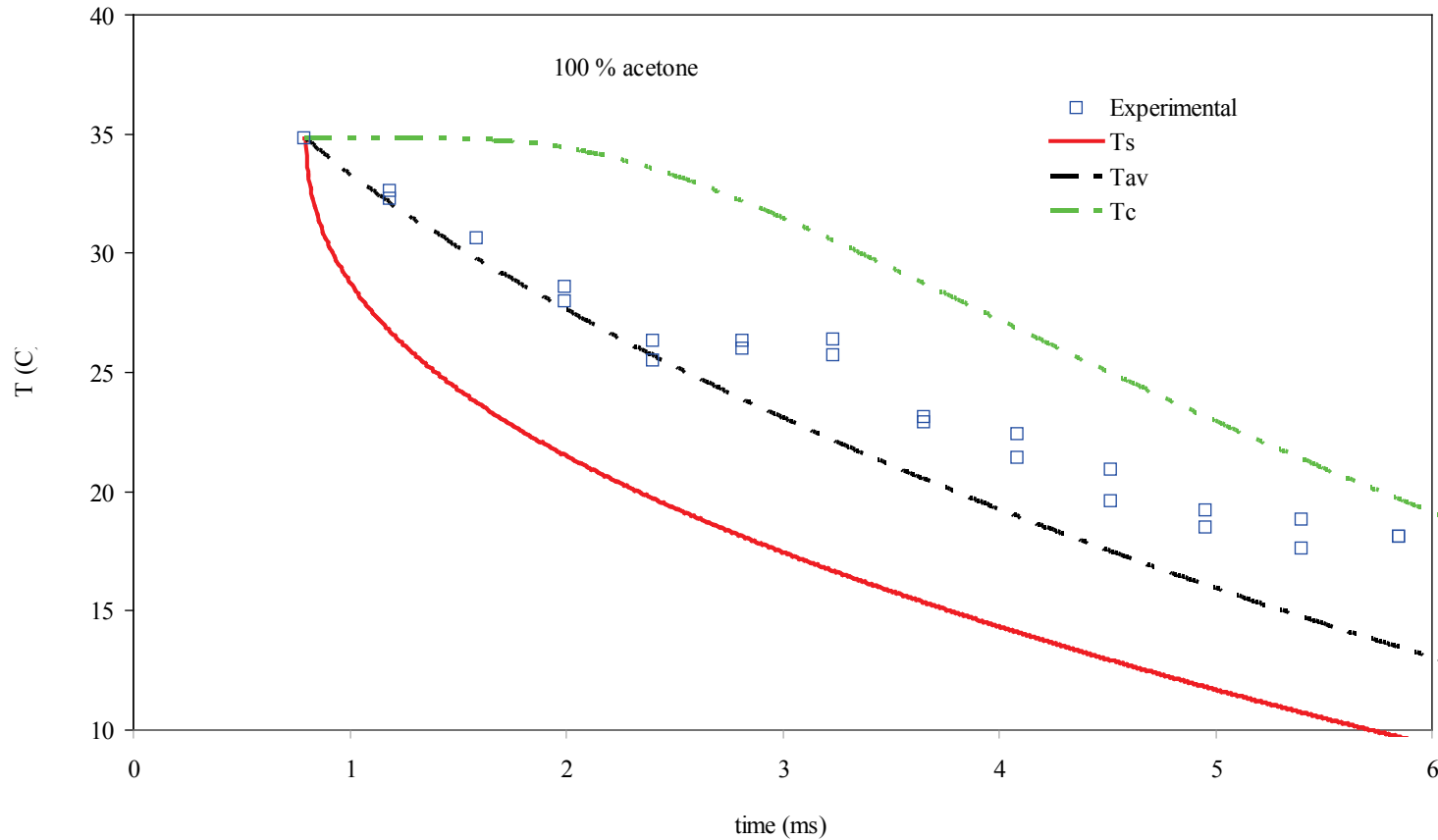
$$\chi_{li} = \frac{Y_{li}/M_i}{\sum_i (Y_{li}/M_i)}$$

$$\chi_{i,g,S} P_a = \chi_{i,l,S} \gamma_i P_{\text{sat}}(T_S)$$

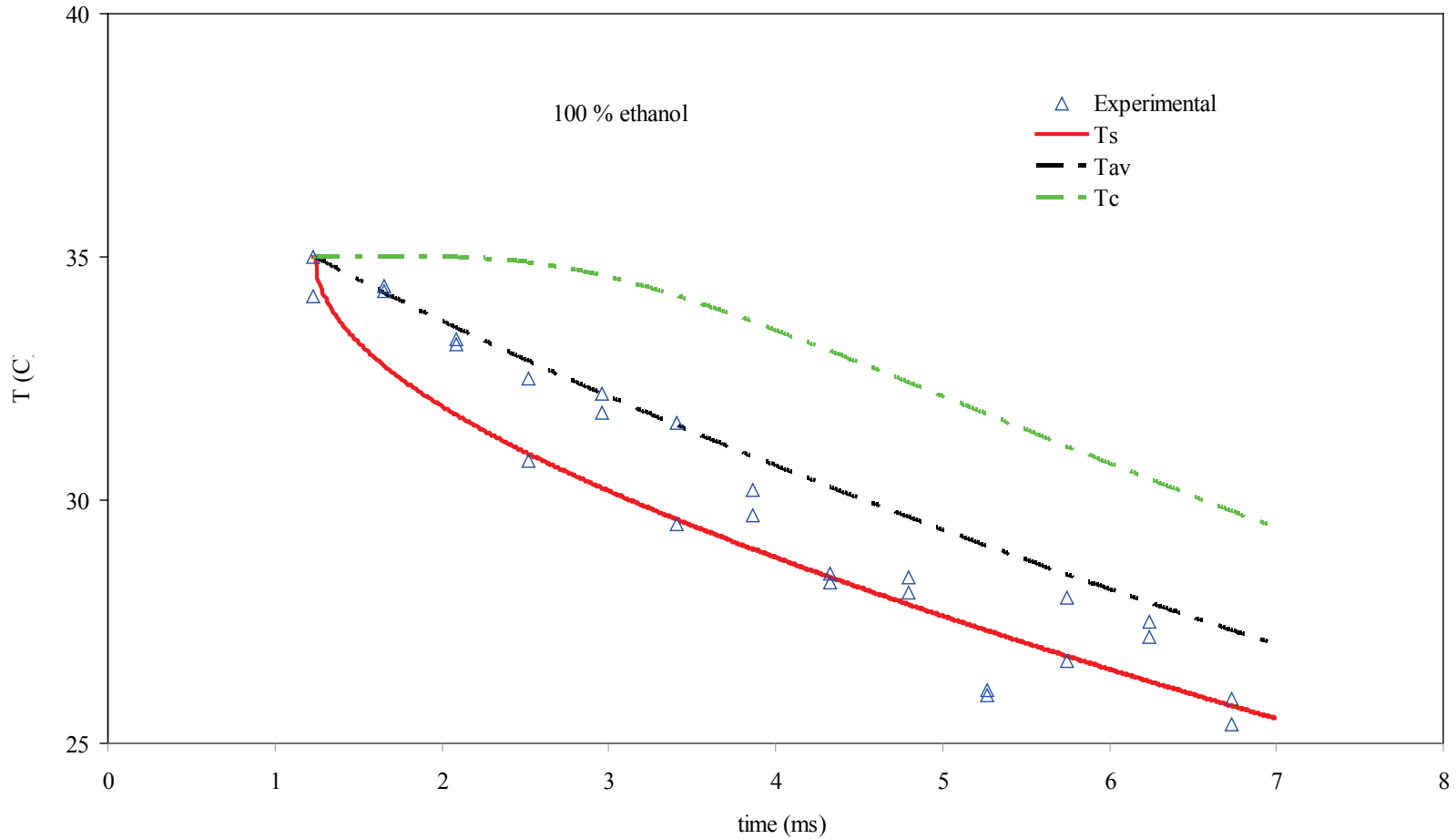
- where γ_i is the activity coefficient and as a first step we put it equal to 1 (ideal mixture).

$$Y_{vi} = \frac{\chi_{vi} M_i}{\sum_i (\chi_{vi} M_i) + \chi_{\text{air}} M_{\text{air}}}$$

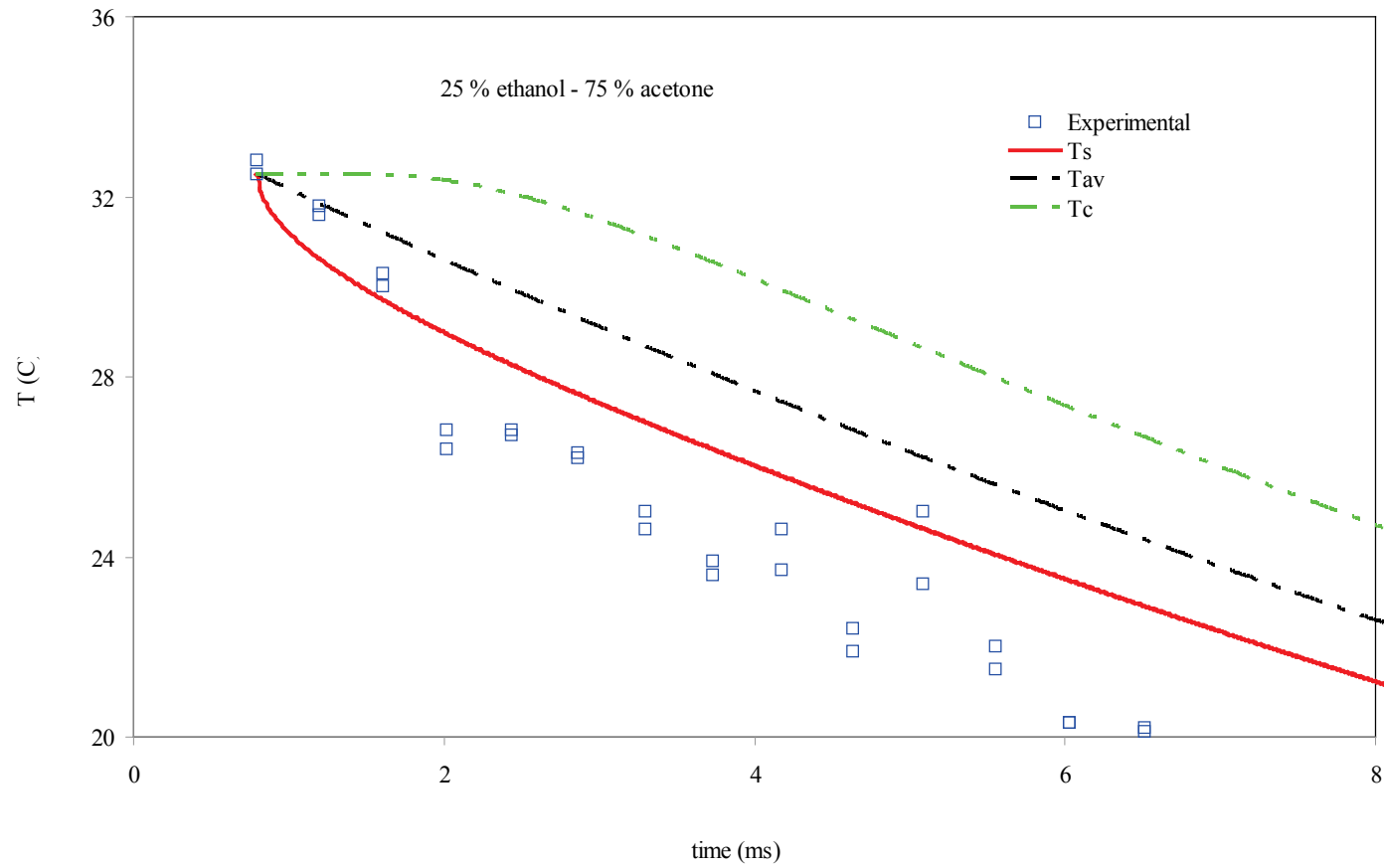
Preliminary results



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Thank you for your attention

Your questions are more than welcome

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